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ON THE OSCILLATION OF A CLASS OF DAMPED FRACTIONAL DIFFERENTIAL EQUATIONS

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Abstract. Using Riccati type transformations, the authors establish some new oscillation criteria for the fractional differential equation

$$\left(D_{0+}^{1+\alpha}y\right)(t) + p(t)\left(D_{0+}^{\alpha}y\right)(t) + q(t)f(G(t)) = 0, \quad t > 0,$$

where $D_{0+}^{\alpha} y$ is the Riemann-Liouville fractional derivative of order α of y, $G(t) = \int_{0}^{t} (t-s)^{-\alpha} y(s) ds$, and $\alpha \in (0, 1)$. Examples are provided to illustrate the relevance of the results.

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1. INTRODUCTION

In this paper we are concerned with the oscillatory nature of the fractional differential equation with damping

$$\left(D_{0+}^{1+\alpha}y\right)(t) + p(t)\left(D_{0+}^{\alpha}y\right)(t) + q(t)f(G(t)) = 0, \qquad t \ge t_0 > 0, \tag{1.1}$$

where $\alpha \in (0, 1)$ is a constant, $D_{0+}^{\alpha} y$ is the α -th Riemann-Liouville fractional derivative of y, and $G(t) = \int_{0}^{t} (t-s)^{-\alpha} y(s) ds$.

We assume throughout this paper that the following conditions hold:

(C1) $p: [t_0, \infty) \to R$ is a continuous function with p(t) < 0;

(C2) $q : [t_0, \infty) \to R$ is a continuous function with $q(t) \ge 0$;

(C3) $f: R \to R$ is a continuous function with uf(u) > 0 for $u \neq 0$ and there exists a constant K > 0 such that $\frac{f(u)}{u} \ge K$ for all $u \neq 0$. A nontrivial function y(t) is said to be a solution of equation (1.1) if $y(t) \in$

 $C((0,\infty), R), \int_{0}^{t} (t-s)^{-\alpha} y(s) ds \in C^{1}((0,\infty), R), D_{0+}^{\alpha} y \in C^{1}((0,\infty), R), \text{ and sat-}$ isfies (1.1) on $(0,\infty)$. Our attention is restricted to those solutions of (1.1) which © 2016 Miskolc University Press

exist on $(0, \infty)$ and satisfy $\sup\{|y(t)| : t \ge t_*\} > 0$ for any $t_* > 0$. Such a solution of equation (1.1) is said to be oscillatory if it is neither eventually positive nor eventually negative, and it is nonoscillatory otherwise. Equation (1.1) is said to be oscillatory if all its solutions are oscillatory.

Differential equations of fractional orders appear more and more frequently in various research areas and applications in diverse fields of science and engineering. Readers are referred to the monographs [2, 7, 8, 10] for the theory and applications of fractional calculus.

The problem of determining the oscillation of solutions of various equations like ordinary differential equations, difference equations, and functional differential equations has been a very active area of research in the last few decades, for example, see [1,5,11,12] and the references therein.

However, to the best of our knowledge, only a few papers considered the oscillatory behavior of fractional differential equations. We refer the reader to [3,6,9,13] and the references therein for problems involving Riemann-Liouville fractional derivatives; for the problems involving Caputo fractional derivatives, the reader is referred to [4,6] and the included references.

The motivation for the present work has been inspired basically by the papers of Chen [3], Prakash et.al [9], and the cited papers in the references. By using Riccati type transformations and the integral averaging technique, we establish some new sufficient conditions which guarantee the oscillation of solutions of equation (1.1). It is therefore hoped that the present paper will contribute to the studies in oscillatory behavior of solutions of fractional differential equations with damping. Finally, some examples are given to illustrate our results.

2. PRELIMINARIES

In this section, we introduce some basic notations, definitions, and preliminary facts from fractional calculus; see [2, 7, 8, 10] for additional details.

Definition 1. The Riemann-Liouville fractional integral of order $\beta > 0$ for a function $y : (0, \infty) \rightarrow R$ is defined by

$$\left(I_{0+}^{\beta}y\right)(t) := \frac{1}{\Gamma(\beta)} \int_{0}^{t} (t-s)^{\beta-1} y(s) ds, \qquad t > 0,$$
(2.1)

provided the right side exists pointwise on $(0, \infty)$, where Γ is the Gamma function.

Definition 2. The Riemann-Liouville fractional derivative of order $\beta > 0$ for a function $y : (0, \infty) \rightarrow R$ is defined by

$$\left(D_{0+}^{\beta}y\right)(t) := \frac{1}{\Gamma(n-\beta)} \frac{d^n}{dt^n} \int_0^t (t-s)^{n-\beta-1} y(s) ds, \qquad t > 0, \qquad (2.2)$$

provided the righ-hand side exists, where $n = [\beta] + 1$ and $[\beta]$ denotes the integer part of β .

To obtain our results in this paper, we need the following two lemmas.

Lemma 1. Let y be solution of (1.1), and

$$G(t) := \int_{0}^{t} (t-s)^{-\alpha} y(s) ds \quad for \ \alpha \in (0,1), \quad t > 0.$$
 (2.3)

Then

$$G'(t) = \Gamma(1-\alpha) \left(D_{0+}^{\alpha} y \right)(t).$$
(2.4)

Proof. Differentiating (2.3) and using (2.2), we have, for $\alpha \in (0, 1)$ and t > 0,

$$G'(t) = \frac{d}{dt} \int_0^t (t-s)^{-\alpha} y(s) ds$$
$$= \Gamma(1-\alpha) \frac{1}{\Gamma(1-\alpha)} \frac{d}{dt} \int_0^t (t-s)^{-\alpha} y(s) ds$$
$$= \Gamma(1-\alpha) \left(D_{0+}^{\alpha} y\right)(t).$$

This completes the proof proof of Lemma 1.

Lemma 2. Let $\alpha \in (0, 1)$, and t > 0. If y is a solution of (1.1), then $(D_{0+}^{1+\alpha}y)(t) = (D_{0+}^{\alpha}y)'(t)$.

Proof. From (2.2) and $n = [\beta] + 1 = [1 + \alpha] + 1 = [\alpha] + 1 + 1 = 2$, we obtain

$$(D_{0+}^{1+\alpha}y)(t) = \frac{1}{\Gamma(2-\alpha-1)} \frac{d^2}{dt^2} \int_0^t (t-s)^{2-\alpha-1-1} y(s) ds$$
$$= \frac{1}{\Gamma(1-\alpha)} \frac{d^2}{dt^2} \int_0^t (t-s)^{-\alpha} y(s) ds$$

$$= \frac{d}{dt} \left(\frac{1}{\Gamma(1-\alpha)} \frac{d}{dt} \int_{0}^{t} (t-s)^{-\alpha} y(s) ds \right)$$
$$= \frac{d}{dt} \left(D_{0+}^{\alpha} y \right)(t) = \left(D_{0+}^{\alpha} y \right)'(t).$$

This completes the proof of Lemma 2.

3. MAIN RESULTS

Theorem 1. Suppose that (C1)-(C3) hold. If

$$\lim_{t \to \infty} \left[\frac{1}{4\Gamma(1-\alpha)} \int_{t_0}^t \left(4\Gamma(1-\alpha) Kq(s) - p^2(s) \right) ds \right] = \infty, \tag{3.1}$$

then any solution of the equation (1.1) is oscillatory.

Proof of Theorem 1. Let y(t) be a nonoscilatory solution of equation (1.1). Without loss of generality, we may assume that y is an eventually positive solution of (1.1). Then, there exists $t_1 \in [t_0, \infty)$ such that y(t) > 0 and G(t) > 0 for $t \ge t_1$. Define the function w by

$$w(t) = -\frac{(D_{0+}^{\alpha}y)(t)}{G(t)}, \quad \text{for } t \ge t_1.$$

Then w(t) is well defined and satisfies the inequality

$$w'(t) = -\left[\frac{\left(\left(D_{0+}^{\alpha}y\right)(t)\right)'G(t) - \left(D_{0+}^{\alpha}y\right)(t)G'(t)\right)}{G^{2}(t)}\right]$$

= $-\left[\frac{\left(D_{0+}^{1+\alpha}y\right)(t)}{G(t)} - \frac{\Gamma(1-\alpha)\left(\left(D_{0+}^{\alpha}y\right)(t)\right)^{2}}{G^{2}(t)}\right]$
= $p(t)\frac{\left(D_{0+}^{\alpha}y\right)(t)}{G(t)} + q(t)\frac{f(G(t))}{G(t)} + \Gamma(1-\alpha)\left(\frac{\left(D_{0+}^{\alpha}y\right)(t)}{G(t)}\right)^{2}$
 $\geq -p(t)w(t) + Kq(t) + \Gamma(1-\alpha)w^{2}(t).$ (3.2)

Integrating (3.2) from t_1 to t, we obtain

$$w(t) \ge w(t_1) + \int_{t_1}^t \left[\Gamma(1-\alpha) w^2(s) - p(s)w(s) + Kq(s) \right] ds$$

= $w(t_1) + \Gamma(1-\alpha) \int_{t_1}^t \left(w(s) - \frac{p(s)}{2\Gamma(1-\alpha)} \right)^2 ds$

650

$$+\frac{1}{4\Gamma(1-\alpha)}\int_{t_1}^t \left(4\Gamma(1-\alpha)Kq(s)-p^2(s)\right)ds.$$
(3.3)

In view of (3.1), there exists $t_2 \ge t_1$, such that

$$w(t) > \Gamma(1-\alpha) \int_{t_2}^t \left(w(s) - \frac{p(s)}{2\Gamma(1-\alpha)} \right)^2 ds \quad \text{for } t \ge t_2.$$

If we set

$$H(t) = \Gamma(1-\alpha) \int_{t_2}^t \left(w(s) - \frac{p(s)}{2\Gamma(1-\alpha)} \right)^2 ds, \qquad (3.4)$$

then we have w(t) > H(t) > 0 for $t \ge t_2$.

From (3.4) and using the fact that p(t) < 0, we can easily see that

$$H'(t) = \Gamma(1-\alpha) \left(w(t) - \frac{p(t)}{2\Gamma(1-\alpha)} \right)^2 > \Gamma(1-\alpha) w^2(t) > \Gamma(1-\alpha) H^2(t),$$

and so we obtain

$$\Gamma(1-\alpha) < \frac{H'(t)}{H^2(t)}$$
 for $t \ge t_2$.

Integrating both sides of this inequality from t_2 to t, we obtain that

$$\Gamma(1-\alpha)\int_{t_2}^{t} ds < \frac{1}{H(t_2)} - \frac{1}{H(t)} < \frac{1}{H(t_2)}.$$

Letting $t \to \infty$ in the above inequality gives

$$\lim_{t\to\infty}\Gamma(1-\alpha)\int\limits_{t_2}^t ds < \frac{1}{H(t_2)},$$

but this is not correct. This completes the proof of the theorem.

Theorem 2. Let conditions (C1)–(C3) hold. Assume that there exists a positive function $g \in C^1[t_0, \infty)$ such that

$$\lim_{t \to \infty} \int_{t_0}^t \frac{\Gamma(1-\alpha)}{g(s)} ds = \infty$$
(3.5)

and

$$\lim_{t \to \infty} \left[-\frac{1}{4\sqrt{\Gamma(1-\alpha)}} \int_{t_0}^t \Psi(s) ds + \frac{g'(s)}{2\sqrt{\Gamma(1-\alpha)}} \right] = \infty,$$
(3.6)

where

$$\Psi(s) = p^2(s)g(s) + \frac{(g'(s))^2}{g(s)} - 2p(s)g'(s) - 4K\Gamma(1-\alpha)g(s)q(s),$$

then any solution of the equation (1.1) is oscillatory.

Proof of Theorem 2. Let y(t) be a nonoscilatory solution of equation (1.1). Without loss of generality we may assume that y is an eventually positive solution of (1.1). Then, there exists $t_1 \in [t_0, \infty)$ such that y(t) > 0 and G(t) > 0 for $t \ge t_1$. Define the function w by

$$w(t) = -g(t)\frac{(D_{0+}^{\alpha}y)(t)}{G(t)} \quad \text{for } t \ge t_1.$$
(3.7)

In view of conditions (C1)-(C3), from equation (1.1) we obtain

$$w'(t) = -g'(t) \frac{(D_{0+}^{\alpha} y)(t)}{G(t)} - g(t) \left[\frac{(D_{0+}^{\alpha} y)(t)}{G(t)} \right]'$$

$$= \frac{g'(t)}{g(t)} w(t) - g(t) \left[\frac{((D_{0+}^{\alpha} y)(t))'G(t) - (D_{0+}^{\alpha} y)(t)G'(t)}{G^{2}(t)} \right]$$

$$= \frac{g'(t)}{g(t)} w(t) - g(t) \left[\frac{(D_{0+}^{1+\alpha} y)(t)}{G(t)} - \frac{\Gamma(1-\alpha)\left((D_{0+}^{\alpha} y)(t)\right)^{2}}{G^{2}(t)} \right]$$

$$= \frac{g'(t)}{g(t)} w(t) + g(t) p(t) \frac{(D_{0+}^{\alpha} y)(t)}{G(t)} + g(t)q(t) \frac{f(G(t))}{G(t)}$$

$$+ \Gamma(1-\alpha)g(t) \left(\frac{(D_{0+}^{\alpha} y)(t)}{G(t)} \right)^{2}$$

$$\geq \frac{1}{g(t)} \left[\Gamma(1-\alpha)w^{2}(t) - g(t)p(t)w(t) + g'(t)w(t) \right] + Kg(t)q(t).$$
(3.8)

Taking

$$H(t) = \sqrt{\Gamma(1-\alpha)}w(t) + \frac{1}{2\sqrt{\Gamma(1-\alpha)}}g'(t),$$
 (3.9)

we get from (3.8) that

$$w'(t) \ge \frac{1}{g(t)} \left[\left(H(t) - \frac{p(t)g(t)}{2\sqrt{\Gamma(1-\alpha)}} \right)^2 - \left(\frac{p(t)g(t)}{2\sqrt{\Gamma(1-\alpha)}} \right)^2 - \frac{(g'(t))^2}{4\Gamma(1-\alpha)} + \frac{p(t)g(t)g'(t)}{2\Gamma(1-\alpha)} \right] + Kg(t)q(t)$$

652

ON THE OSCILLATION OF A CLASS OF DAMPED FRACTIONAL DIFFERENTIAL EQUATIONS 653

$$= \frac{1}{g(t)} \left(H(t) - \frac{p(t)g(t)}{2\sqrt{\Gamma(1-\alpha)}} \right)^2 - \frac{1}{4\Gamma(1-\alpha)} \left[p^2(t)g(t) + \frac{(g'(t))^2}{g(t)} - 2p(t)g'(t) - 4K\Gamma(1-\alpha)g(t)q(t) \right].$$

Integrating both of sides of the above equation from t_1 to t, we obtain

$$w(t) \ge w(t_1) + \int_{t_1}^t \frac{1}{g(s)} \left(H(s) - \frac{p(s)g(s)}{2\sqrt{\Gamma(1-\alpha)}} \right)^2 ds$$
(3.10)

$$-\frac{1}{4\Gamma(1-\alpha)}\int_{t_1}^t \left[p^2(s)g(s) + \frac{(g'(s))^2}{g(s)} - 2p(s)g'(s) - 4K\Gamma(1-\alpha)g(s)q(s)\right]ds.$$

Using (3.9) in (3.10), we obtain

$$\begin{split} H(t) &\geq \sqrt{\Gamma(1-\alpha)}w(t_1) + \sqrt{\Gamma(1-\alpha)} \int_{t_1}^t \frac{1}{g(s)} \left(H(s) - \frac{p(s)g(s)}{2\sqrt{\Gamma(1-\alpha)}} \right)^2 ds \\ &- \frac{1}{4\sqrt{\Gamma(1-\alpha)}} \int_{t_1}^t \left[p^2(s)g(s) + \frac{\left(g'(s)\right)^2}{g(s)} - 2p(s)g'(s) - 4K\Gamma(1-\alpha)g(s)q(s) \right] ds \\ &+ \frac{g'(t)}{2\sqrt{\Gamma(1-\alpha)}}. \end{split}$$

Now, from (3.6) there exists $t_2 \ge t_1$ such that

$$H(t) > \sqrt{\Gamma(1-\alpha)} \int_{t_1}^t \frac{1}{g(s)} \left(H(s) - \frac{p(s)g(s)}{2\sqrt{\Gamma(1-\alpha)}} \right)^2 ds.$$

Let

$$Q(t) = \sqrt{\Gamma(1-\alpha)} \int_{t_1}^t \frac{1}{g(s)} \left(H(s) - \frac{p(s)g(s)}{2\sqrt{\Gamma(1-\alpha)}} \right)^2 ds.$$
(3.11)

Then, using the fact that p(t) < 0, we have H(t) > Q(t) > 0. Differentiating (3.11) gives

$$Q'(t) = \sqrt{\Gamma(1-\alpha)} \frac{1}{g(t)} \left(H(t) - \frac{p(t)g(t)}{2\sqrt{\Gamma(1-\alpha)}} \right)^2$$
$$\geq \sqrt{\Gamma(1-\alpha)} \frac{1}{g(t)} \left(Q(t) - \frac{p(t)g(t)}{2\sqrt{\Gamma(1-\alpha)}} \right)^2$$

ERCAN TUNÇ AND OSMAN TUNÇ

$$>\sqrt{\Gamma(1-\alpha)}\frac{1}{g(t)}Q^2(t),$$

and so we have

$$\frac{\sqrt{\Gamma(1-\alpha)}}{g(t)} < \frac{Q'(t)}{Q^2(t)}$$

Integrating both of sides of this inequality from t_2 to t, we obtain

$$\int_{t_2}^{t} \frac{\Gamma(1-\alpha)}{g(s)} ds < \frac{1}{Q(t_2)} - \frac{1}{Q(t)} < \frac{1}{Q(t_2)},$$

which contradicts (3.5) and completes the proof of the theorem.

Example 1. Consider the fractional differential equation, for t > 0,

$$\left(D_{0+}^{1+\alpha}y\right)(t) - \frac{1}{t^2}\left(D_{0+}^{\alpha}y\right)(t) + \left(1 + \frac{1}{t^6}\right)\int_0^t (t-s)^{-\alpha}y(s)ds = 0, \quad (3.12)$$

with $\alpha = 1/2$. Here, $p(t) = -1/t^2$, $q(t) = 1 + 1/t^6$, and f(u) = u. Then we have K = 1, and $\Gamma(1-\alpha) = \Gamma(1/2) = \sqrt{\pi}$. Let $t_0 = 1$, thus condition (3.1) becomes

$$\lim_{t \to \infty} \left[\frac{1}{4\Gamma(1-\alpha)} \int_{t_0}^t \left(4\Gamma(1-\alpha) Kq(s) - p^2(s) \right) ds \right]$$
$$= \lim_{t \to \infty} \left[\frac{1}{4\sqrt{\pi}} \int_{1}^t \left(4\sqrt{\pi}(1+\frac{1}{s^6}) - \frac{1}{s^4} \right) ds \right]$$
$$= \lim_{t \to \infty} \left[\frac{1}{4\sqrt{\pi}} \left(4\sqrt{\pi}s - \frac{4\sqrt{\pi}}{5s^5} + \frac{1}{3s^3} \right) \Big|_{1}^t \right] = \infty.$$

So every solution of (3.12) is oscillatory by Theorem 1

Example 2. Consider the fractional differential equation

$$(D_{0+}^{1+\alpha}y)(t) - \frac{1}{t} (D_{0+}^{\alpha}y)(t)$$

$$+ \left(1 + \frac{1}{t^2}\right) \frac{1}{\sqrt{\pi}} \exp\left(\int_0^t (t-s)^{-\alpha}y(s)ds\right) \int_0^t (t-s)^{-\alpha}y(s)ds = 0,$$
(3.13)

with $\alpha = 1/2$. Here, p(t) = -1/t, $q(t) = 1 + 1/t^2$, and $f(u) = \frac{1}{\sqrt{\pi}}e^u u$. Then, we have $K = 1/\sqrt{\pi}$, and $\Gamma(1-\alpha) = \Gamma(1/2) = \sqrt{\pi}$. If we take $t_0 = 1$, and g(t) = t,

654

then

$$\lim_{t \to \infty} \int_{t_0}^t \frac{\Gamma(1-\alpha)}{g(s)} ds = \lim_{t \to \infty} \int_{1}^t \frac{\sqrt{\pi}}{s} ds = \infty$$

and condition (3.6) becomes

$$\lim_{t \to \infty} \left[-\frac{1}{4\pi^{1/4}} \int_{1}^{t} \left\{ \frac{1}{s} + \frac{1}{s} + \frac{2}{s} - 4\frac{1}{\sqrt{\pi}} \sqrt{\pi} s \left(1 + \frac{1}{s^2} \right) \right\} ds + \frac{1}{2\pi^{1/4}} \right]$$
$$= \lim_{t \to \infty} \left[-\frac{1}{4\pi^{1/4}} \int_{1}^{t} \{-4s\} ds + \frac{1}{2\pi^{1/4}} \right]$$
$$= \lim_{t \to \infty} \left[\frac{1}{2\pi^{1/4}} s^2 \int_{1}^{t} + \frac{1}{2\pi^{1/4}} \right] = \infty.$$

We see that the hypotheses of Theorem 2 are satisfied and so all solutions of (3.13) are oscillatory.

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